

Difficult instances of the counting problem for 2-quantum-SAT are very atypical

Niel de Beaudrap

CWI, Amsterdam

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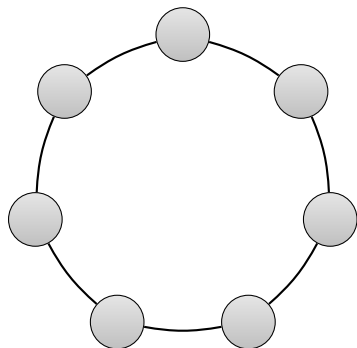
preamble: about frustration-freeness

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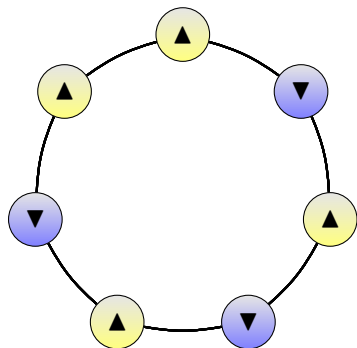
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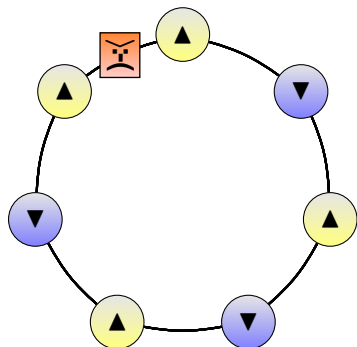
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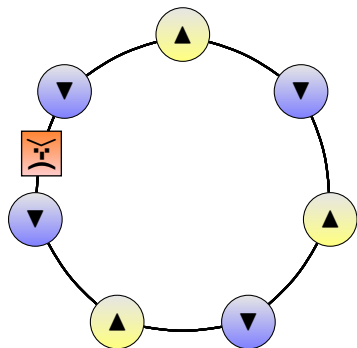
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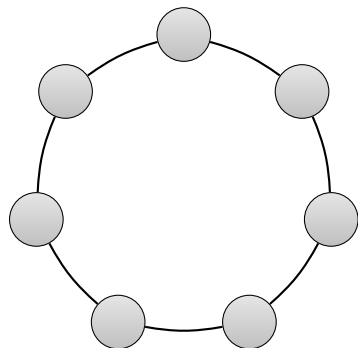
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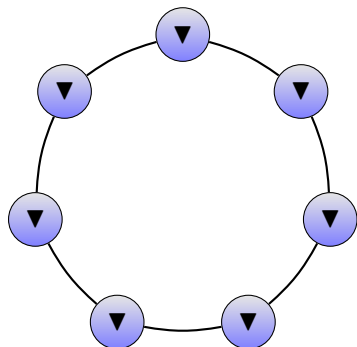
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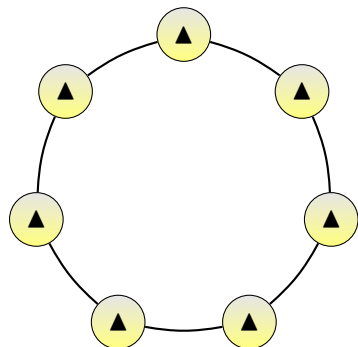
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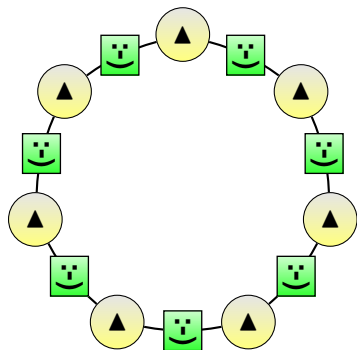
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- (Complexity of **3-SAT/3-QSAT** (3-body constraints): ... very hard.)

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- Generalizes **#2-SAT**: counting solutions to an instance of **2-SAT**

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 - ▶ “entangled constraints” act similarly to a *pair* of product constraints
- More fundamentally: constraints in **multiple local bases**
 - ▶ *Monotonicity* seems to be the source of difficulty of #2-SAT.
 - ▶ The sharp decline in difficulty of #2-QSAT provides evidence for this: monotonicity requires a preferred basis.

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Then as $\mathbf{q} \rightarrow \mathbf{0}$, ‘hard’ instances of **#2-QSAT** become unlikely for any fixed density of interaction terms.

Outline

1 Constructing models of random #2-QSAT

- Erdős–Rényi graphs and percolated lattices
- Random instances of **2-QSAT** on random graphs
- Random frustration-free Hamiltonians on random graphs
- Common features of the interaction graph models

2 Analysis of #2-QSAT on random graphs

- Effective long-range constraints
- Onset of frustration in random two-body Hamiltonians on qubits
- Frozen subsystems in frustration-free models

3 Summary

models of random graphs (part 1)

Prior work on random 2-satisfiability:

- random **2-SAT** with uniformly random boolean constraints *e.g.*, [Chvátal+Reed, 1992]
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These graphs have “as little structure as possible” — while perhaps unphysical, this (and prior work) **motivates** this model.

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Results for square/cubic lattices will have analogues for other lattices.

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according to $\mathbf{q} = (q_1, q_2, \dots)$, where $q_k = \Pr[\alpha = \alpha_k]$
- 3 Repeat until the interaction graph has $m = \gamma n$ edges:

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- **Super-critical graph:** largest component is “complicated” and has size $O(n)$ for $\gamma > \gamma_0$: *interesting* instances.

Outline

1

Constructing models of random #2-QSAT

- Erdős–Rényi graphs and percolated lattices
- Random instances of #2-QSAT on random graphs
- Random frustration-free Hamiltonians on random graphs
- Common features of the interaction graph models

2

Analysis of #2-QSAT on random graphs

- Effective long-range constraints
- Onset of frustration in random two-body Hamiltonians on qubits
- Frozen subsystems in frustration-free models

3

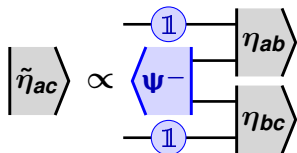
Summary

effective next-nearest-neighbor constraints

- We describe constraints $h_{ab} = |\eta_{ab}\rangle\langle\eta_{ab}|$ in terms of the state $|\eta_{ab}\rangle$ in its support.

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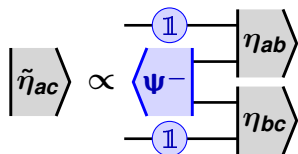
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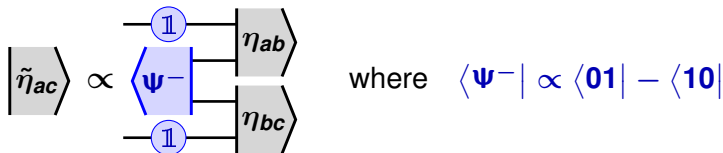


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- For constraints $|\eta\rangle = |\alpha\rangle|\alpha'\rangle$ whose factors are distributed *i.i.d.* according to \mathbf{q} , $|\tilde{\eta}_{ac}\rangle \neq \mathbf{0}$ happens with probability

$$Q := \sum_{k \geq 1} q_k(1 - q_k) = 1 - \sum_{k \geq 1} q_k^2 = 1 - \|\mathbf{q}\|_2^2.$$

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Ansatz: long range constraints in a Hamiltonian with density $\gamma > 0$ act like connectivity in a random graph with density γQ

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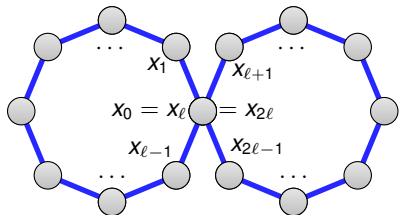
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- **Claim:** Random Hamiltonians on qubits have a phase transition from **unfrustration** \rightarrow **frustration** consistent with this “edge attenuation” ansatz

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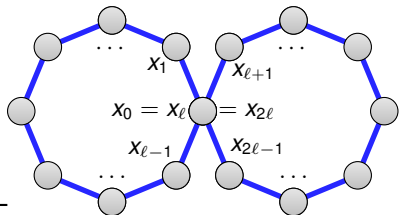
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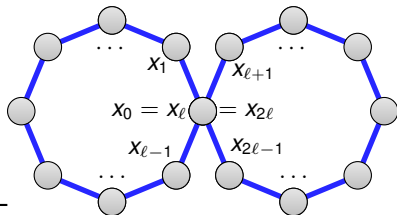
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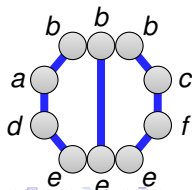
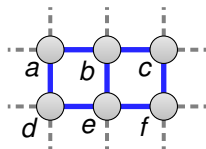
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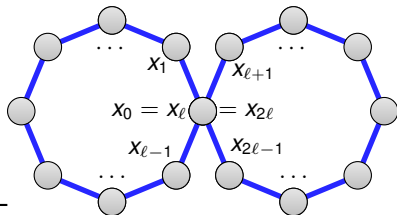
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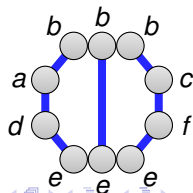
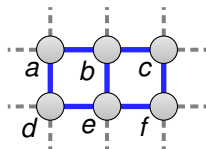
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⇒ super-critical random Hamiltonians on percolated lattices are almost surely frustrated.



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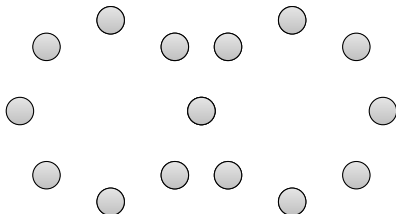
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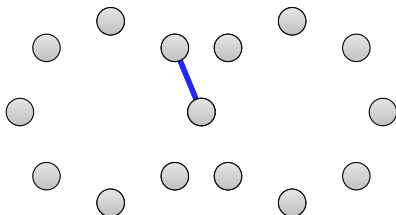
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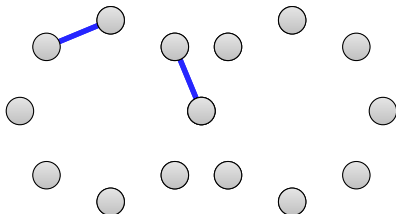
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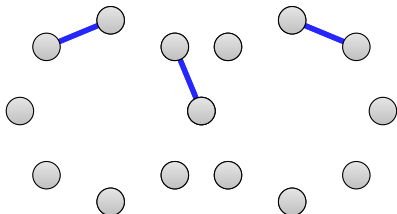
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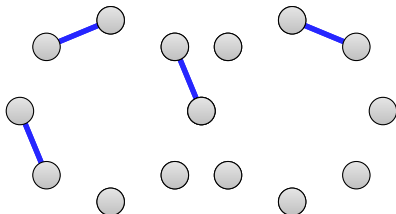
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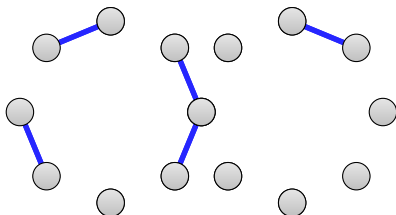
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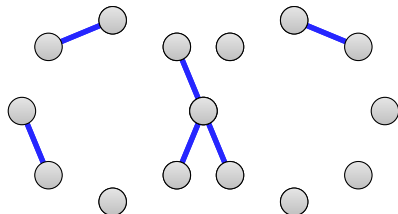
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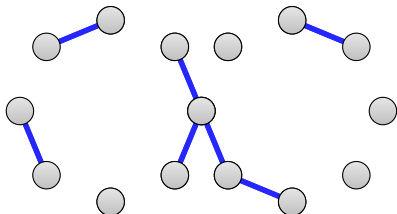
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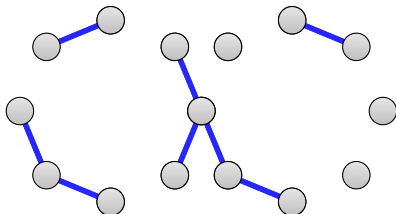
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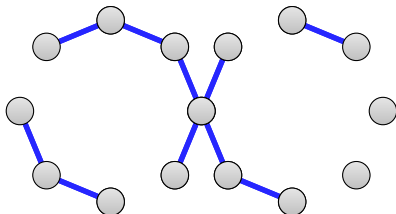
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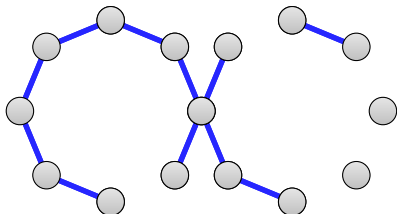
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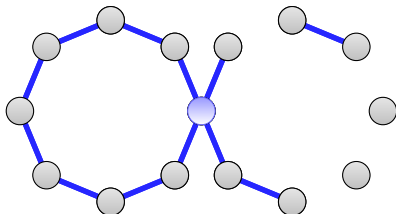
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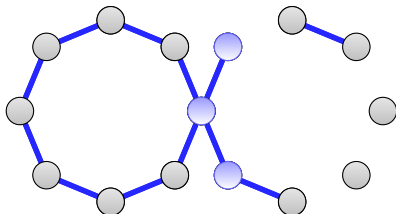
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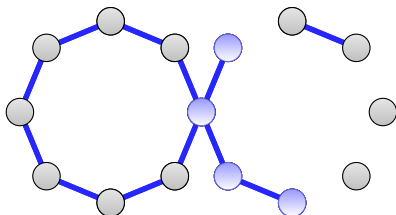
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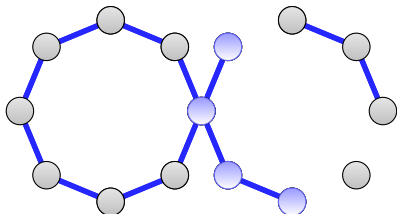
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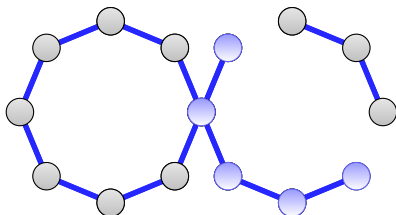
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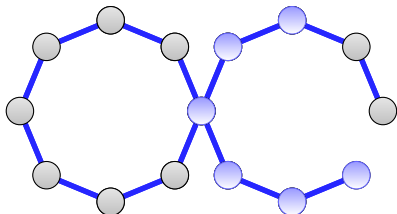
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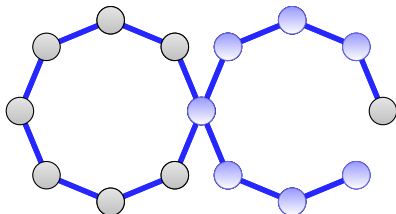
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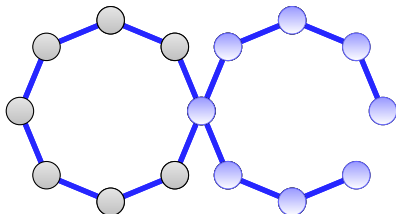
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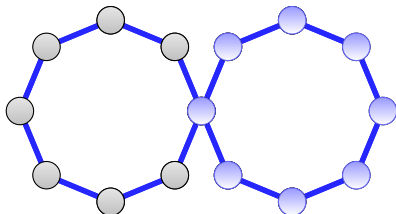
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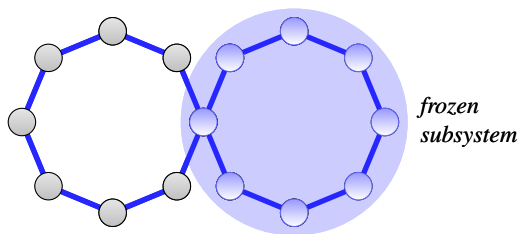
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- A dense enough frustration-free Hamiltonian will almost surely have a large **frozen subsystem**, consisting of qubits whose states are fixed by the Hamiltonian

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The **frozen subsystems** form a random subgraph of the random interaction graph.

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- **Ansatz.** The frozen graph acts like a random graph with some missing components, and edge-density “similar” to γQ .

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 - ▶ Simulate growth of the frozen subgraph within G by adding edges *i.i.d.* with probability Q_∞ .
 - ▶ Any frozen subsystem of size $\omega(\log n)$ can only exist within a “giant” frozen component, whose growth is then bounded from below.

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 - ▶ As $q \rightarrow 0$: eventually no room for a “potentially difficult” regime
- **Percolated lattices:** For each $d \in \{2, 3\}$ there is a percolation constant $0 < p_{\text{fin}} < 1$ such that if $Q_\infty > p_{\text{fin}}$, the frozen subgraph almost surely decouples Hamiltonians of *any* edge-density into tractible pieces

($p_{\text{fin}} = \frac{1}{2}$ for square lattices; $p_{\text{fin}} \geq p_c \approx 0.24881$ for cubic lattices)

Outline

1

Constructing models of random #2-QSAT

- Erdős–Rényi graphs and percolated lattices
- Random instances of #2-QSAT on random graphs
- Random frustration-free Hamiltonians on random graphs
- Common features of the interaction graph models

2

Analysis of #2-QSAT on random graphs

- Effective long-range constraints
- Onset of frustration in random two-body Hamiltonians on qubits
- Frozen subsystems in frustration-free models

3

Summary

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